# ON SOME INEQUALITIES FOR MEANS AND THE SECOND GAUTSCHI-KERSHAW'S INEQUALITY 

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Abstract. We establish some inequalities for means, and we present new bounds for the second Gautschi-Kershaw's inequality.

## 1. Introduction

Let $r, s \in \mathbb{R}$ and let $a, b>0$. The Stolarsky mean $E_{r, s}(a, b)$ of order $(r, s)$ of $a$ and $b$ with $a \neq b$ are defined as

$$
E_{r, s}(a, b)= \begin{cases}\left(\frac{r}{s} \cdot \frac{b^{s}-a^{s}}{b^{r}-a^{r}}\right)^{1 /(s-r)}, & r s(r-s) \neq 0  \tag{1}\\ \exp \left(-\frac{1}{r}+\frac{a^{r} \ln a-b^{r} \ln b}{a^{r}-b^{r}}\right), & r=s \neq 0 \\ \left(\frac{1}{r} \cdot \frac{b^{r}-a^{r}}{\ln b-\ln a}\right)^{1 / r}, & r \neq 0, s=0 \\ \sqrt{a b}, & r=s=0\end{cases}
$$

with $E_{r, s}(a, a)=a$ (see [51, 52]), while the Gini mean $G_{r, s}(a, b)$ of order $(r, s)$ of $a$ and $b$ are defined in 18 as

$$
G_{r, s}(a, b)= \begin{cases}\left(\frac{a^{s}+b^{s}}{a^{r}+b^{r}}\right)^{1 /(s-r)}, & r \neq s  \tag{2}\\ \exp \left(\frac{a^{r} \ln a+b^{r} \ln b}{a^{r}+b^{r}}\right), & r=s \neq 0 \\ \sqrt{a b}, & r=s=0\end{cases}
$$

K. B. Stolarsky [51], Leach and Sholander [27] showed that $E_{r, s}(a, b)$ are for $a \neq b$ strictly increasing with both $r$ and $s$. For $a \neq b, G_{r, s}(a, b)$ are also strictly increasing with both $r$ and $s$, see [36, 42]. Leach and Sholander [29] and Páles [38] solved the problem of comparison of Stolarsky mean. The problem of comparison of Gini mean was completely solved by Páles [39] (see also the paper by P. Czinder and Zs. Ples [12]). A problem of comparability of Gini and Stolarsky means was addressed by Neuman and Páles [35]. Minkowski-type inequality for Stolarsky and Gini means can be found in [11, 31, 32].

[^0]Since $E_{r, s}(a, b)$ are for $a \neq b$ strictly increasing with both $r$ and $s$, for the particular choices of the parameters $r$ and $s$, we obtain the following chain of inequalities:

$$
E_{-2,-1}(a, b)<E_{0,0}(a, b)<E_{1,0}(a, b)<E_{1,1}(a, b)<E_{2,1}(a, b) \quad \text { for } \quad a \neq b,
$$

that is,

$$
H(a, b)<G(a, b)<L(a, b)<I(a, b)<A(a, b) \quad \text { for } \quad a \neq b,
$$

where $H, G, L, I$ and $A$ are the harmonic, geometric, logarithmic, identric and arithmetic means, respectively.

It is worth mentioning that

$$
G_{0, r}(a, b)=E_{r, 2 r}(a, b)=M_{r}(a, b)= \begin{cases}\left(\frac{a^{r}+b^{r}}{2}\right)^{1 / r}, & r \neq 0 \\ \sqrt{a b}, & r=0\end{cases}
$$

Thus the classes of Gini and Stolarsky means contain both the power means. Alzer and Ruscheweyh [2] have proven that the joint elements in the classes of Gini and Stolarsky means are exactly the power means.

Let $r, s \in \mathbb{R}$ and $a, b>0$, the generalized Muirhead mean $\sum_{r, s}(a, b)$ of $a$ and $b$ is defined by (see, for instance, [5, p. 333] or [4])

$$
\sum_{r, s}(a, b)=\left(\frac{a^{r} b^{s}+a^{s} b^{r}}{2}\right)^{1 /(r+s)}, r+s \neq 0
$$

T. Trif [53] investigated the monotonicity of $\sum_{r, s}(a, b)$ with respect to $r$ or $s$. Likewise, a comparison theorem and a Minkowski-type inequality involving the generalized Muirhead mean $\sum_{r, s}(a, b)$ are established.

In the special case when $r+s=1$, i.e., $r=\alpha, s=1-\alpha$, the Muirhead (or symmetric) mean is obtained:

$$
\sum_{\alpha, 1-\alpha}(a, b)=\frac{a^{\alpha} b^{1-\alpha}+a^{1-\alpha} b^{\alpha}}{2}
$$

Following A. O. Pittenger [41], we write symmetric mean into the from

$$
S_{\delta}(a, b)=\frac{a^{\frac{1+\sqrt{\delta}}{2}} b^{\frac{1-\sqrt{\delta}}{2}}+a^{\frac{1-\sqrt{\delta}}{2}} b^{\frac{1+\sqrt{\delta}}{2}}}{2}
$$

It is shown in [22] that $S_{\delta}$ is increasing in $\delta$ and that for $a \neq b$

$$
M_{0}(a, b)<S_{\delta}(a, b)<M_{1}(a, b)
$$

provided $0<\delta<1$.
The following result is known

$$
\begin{equation*}
S_{1 / 3}(a, b)<L(a, b)<M_{1 / 3}(a, b) \tag{3}
\end{equation*}
$$

The first inequality in (3) has been established by A. O. Pittenger [41, while the second one was proven in [30, 41. The first inequality in (3) improves a result of B. C. Carlson [9], who proved $S_{1 / 4}(a, b)<L(a, b)$.
E. Neuman [34] established the integral representation

$$
\begin{equation*}
L(a, b)=\int_{0}^{1} a^{t} b^{1-t} d t \tag{4}
\end{equation*}
$$

By applying the Gausss quadrature formula with two knots (see [13, pp. 343-344], [14, p. 36])

$$
\begin{equation*}
\int_{0}^{1} f(t) d t=\frac{1}{2} f\left(\frac{1}{2}+\frac{1}{2 \sqrt{3}}\right)+\frac{1}{2} f\left(\frac{1}{2}-\frac{1}{2 \sqrt{3}}\right)+\frac{1}{4320} f^{(4)}(\xi), \quad 0<\xi<1 \tag{5}
\end{equation*}
$$

to the function $f(t)=a^{t} b^{1-t}$, T. Trif [53] presented a very short proof of the first inequality in (3).

Motivated by the technique of T. Trif, we here present a very short proof of the second inequality in (3). To this aim, by applying Simpson's $\frac{3}{8}$ rule (see [6, 26])

$$
\begin{align*}
\int_{a}^{b} f(t) d t=\frac{b-a}{8}[f(a) & \left.+3 f\left(\frac{2 a+b}{3}\right)+3 f\left(\frac{a+2 b}{3}\right)+f(b)\right]  \tag{6}\\
& -\frac{(b-a)^{5}}{6480} f^{(4)}(\xi) \quad \text { for some } \xi \text { between } a \text { and } b,
\end{align*}
$$

with $a=0, b=1$ and $f(t)=a^{t} b^{1-t}(0 \leq t \leq 1)$, we get

$$
L(a, b)=M_{1 / 3}(a, b)-\frac{1}{6480} a^{\xi} b^{1-\xi}(\ln a-\ln b)^{4} .
$$

This yields the second inequality in (3).
We remark that F. Burk [7] obtained the second inequality in (3) by applying (6) to the function $f(t)=e^{t}$, replacing $a$ and $b$ with $\ln a$ and $\ln b$, respectively.

This paper is organized as follows. In Section 2 we establish some inequalities for means. In Section 3 we presents new bounds, and the series representation for the ratio $\frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x}$, where $\Gamma$ denotes the gamma function.

## 2. Inequalities for means

Theorem 1. Let $a, b>0$ with $a \neq b$, then

$$
\begin{align*}
L & <\frac{1}{3} H+\frac{2}{3} A  \tag{7}\\
\frac{1}{L} & <\frac{1}{3} \frac{1}{H}+\frac{2}{3} \frac{1}{A}  \tag{8}\\
I^{2} & <\frac{1}{3} G^{2}+\frac{2}{3} A^{2} . \tag{9}
\end{align*}
$$

Proof. B. C. Carlson [9] has established the integral representation

$$
\begin{equation*}
L(a, b)=\left[\int_{0}^{1} \frac{1}{t a+(1-t) b} d t\right]^{-1} \tag{10}
\end{equation*}
$$

Use the change of variable $u=t a+(1-t) b$, (10) can be written as

$$
\begin{equation*}
L(a, b)=\left[\frac{1}{b-a} \int_{a}^{b} \frac{1}{u} d u\right]^{-1} \tag{11}
\end{equation*}
$$

By applying (5) to the function $f(t)=\frac{1}{t a+(1-t) b}$, we obtain

$$
\begin{aligned}
\frac{1}{L(a, b)} & =\frac{1}{2}\left[\frac{1}{\left(\frac{1}{2}+\frac{1}{2 \sqrt{3}}\right) a+\left(\frac{1}{2}-\frac{1}{2 \sqrt{3}}\right) b}+\frac{1}{\left(\frac{1}{2}-\frac{1}{2 \sqrt{3}}\right) a+\left(\frac{1}{2}+\frac{1}{2 \sqrt{3}}\right) b}\right] \\
& +\frac{(b-a)^{4}}{180[\xi a+(1-\xi) b]^{5}} \\
= & \frac{1}{\frac{1}{3} H+\frac{2}{3} A}+\frac{(b-a)^{4}}{180[\xi a+(1-\xi) b]^{5}} \quad(0<\xi<1) \\
& >\frac{1}{\frac{1}{3} H+\frac{2}{3} A} .
\end{aligned}
$$

This proves (7).
By applying the composite Simpson rule (see [21])

$$
\begin{align*}
\int_{a}^{b} f(t) d t=\frac{b-a}{6}[f(a) & \left.+4 f\left(\frac{a+b}{2}\right)+f(b)\right]  \tag{12}\\
& -\frac{(b-a)^{5}}{2880} f^{(4)}(\xi) \text { for some } \xi \text { between } a \text { and } b,
\end{align*}
$$

to the function $f(t)=\frac{1}{t}$, we obtain

$$
\frac{1}{L(a, b)}=\frac{1}{6}\left[\frac{1}{b}+\frac{4}{(a+b) / 2}+\frac{1}{a}\right]-\frac{(b-a)^{4}}{120 \xi^{5}}<\frac{1}{3 H}+\frac{2}{3 A}
$$

This proves (8).
It is easy to see that

$$
\begin{equation*}
\ln I(a, b)=\frac{1}{b-a} \int_{a}^{b} \ln t d t=\int_{0}^{1} \ln [t a+(1-t) b] d t \tag{13}
\end{equation*}
$$

By applying (5) to the function $f(t)=\ln [t a+(1-t) b]$, we obtain

$$
\begin{aligned}
\ln I(a, b)= & \frac{1}{2} \ln \left[\left(\frac{1}{2}+\frac{1}{2 \sqrt{3}}\right) a+\left(\frac{1}{2}-\frac{1}{2 \sqrt{3}}\right) b\right] \\
& +\frac{1}{2} \ln \left[\left(\frac{1}{2}-\frac{1}{2 \sqrt{3}}\right) a+\left(\frac{1}{2}+\frac{1}{2 \sqrt{3}}\right) b\right]-\frac{(b-a)^{4}}{720[\xi a+(1-\xi) b]^{4}} \\
= & \frac{1}{2} \ln \left(\frac{1}{6} a^{2}+\frac{2}{3} a b+\frac{1}{6} b^{2}\right)-\frac{(b-a)^{4}}{720[\xi a+(1-\xi) b]^{4}} \\
= & \frac{1}{2} \ln \left(\frac{1}{3} G^{2}+\frac{2}{3} A^{2}\right)-\frac{(b-a)^{4}}{720[\xi a+(1-\xi) b]^{4}} \quad(0<\xi<1) \\
< & \frac{1}{2} \ln \left(\frac{1}{3} G^{2}+\frac{2}{3} A^{2}\right)
\end{aligned}
$$

thus, (9) holds. The proof of Theorem 1 is complete.
Remark 1. The following result is known

$$
\begin{equation*}
\sqrt[3]{A G^{2}}<L<\frac{1}{3} A+\frac{2}{3} G \tag{14}
\end{equation*}
$$

The first inequality in (14) was established by E. B. leach and M. C. Sholander [28], while the second one was proven by B. C. Carlson 9]. The second inequality in (14) can be obtained by applying (12) with $a=0, b=1$ and $f(t)=x^{t} y^{1-t}$.

The inequality

$$
\begin{equation*}
I>\sqrt[3]{G A^{2}} \tag{15}
\end{equation*}
$$

can be concluded by applying (12) to the function $f(t)=\ln t$, see 49, 50. A stronger inequality than (15) is (cf. [49])

$$
\begin{equation*}
I>\frac{1}{3} G+\frac{2}{3} A . \tag{16}
\end{equation*}
$$

## 3. The second Gautschi-Kershaw's inequality

In 1959 W . Gautschi [17] presented the remarkable inequality:

$$
\begin{equation*}
n^{1-s}<\frac{\Gamma(n+1)}{\Gamma(n+s)}<\exp [(1-s) \psi(n+1)], \quad 0<s<1, n=1,2, \ldots \tag{17}
\end{equation*}
$$

where $\psi=\Gamma^{\prime} / \Gamma$ denotes the logarithmic derivative of the gamma function. In 1983 D. Kershaw [24] gave the following closer bounds:

$$
\begin{gather*}
\left(x+\frac{s}{2}\right)^{1-s}<\frac{\Gamma(x+1)}{\Gamma(x+s)}<\left(x-\frac{1}{2}+\sqrt{x+\frac{1}{4}}\right)^{1-s}  \tag{18}\\
\exp [(1-s) \psi(x+\sqrt{s})]<\frac{\Gamma(x+1)}{\Gamma(x+s)}<\exp \left[(1-s) \psi\left(x+\frac{s+1}{2}\right)\right] \tag{19}
\end{gather*}
$$

for real $x>0$ and $0<s<1$. Inequalities (18) and (19) are respectively called the first and the second Gautschi-Kershaw's inequality in the literature.
C. Giordano et al. [19] and B. Palumbo [40] gave a unified treatment and some extensions of Gautschi-Kershaw type inequalities. For each $s>0, x>0$, the inequality (18) is valid for $0<s<1$ or $s>2$, while the reverse inequality is valid for $1<s<2$; and the inequality (19) is valid for $0<s<1$, while the reverse inequality is valid for $s>1$.

The inequality (19) can be written as

$$
\begin{equation*}
\psi(x+\sqrt{s})<\frac{\ln \Gamma(x+1)-\ln \Gamma(x+s)}{1-s}<\psi\left(x+\frac{s+1}{2}\right) . \tag{20}
\end{equation*}
$$

In 2005, D. Kershaw [25] proved that for $s, t>0$,

$$
\begin{equation*}
\psi(x+\sqrt{s t})<\frac{\ln \Gamma(x+t)-\ln \Gamma(x+s)}{t-s}<\psi\left(x+\frac{s+t}{2}\right) . \tag{21}
\end{equation*}
$$

N. Elezović and J. Pec̆arić [16, Lemma 1] proved that for $s, t>0$,

$$
\begin{equation*}
\psi(L(s, t)) \leq \frac{\ln \Gamma(t)-\ln \Gamma(s)}{t-s} \tag{22}
\end{equation*}
$$

N. Batir [3, Theorem 2.7] established an extended form of (22): Let $x$ and $y$ be positive real numbers and $n$ be a positive integer. Then

$$
\begin{align*}
(-1)^{n} \psi^{(n+1)}\left(\frac{x+y}{2}\right)<\frac{(-1)^{n}\left[\psi^{(n)}(x)-\psi^{(n)}(y)\right]}{x} &  \tag{23}\\
& <y \\
& (-1)^{n} \psi^{(n+1)}\left(S_{-(n+1)}(x, y)\right)
\end{align*}
$$

Recently, F. Qi et al. 44] presented closer lower bound:

$$
\begin{equation*}
(-1)^{n} \psi^{(n+1)}(I(x, y))<\frac{(-1)^{n}\left[\psi^{(n)}(x)-\psi^{(n)}(y)\right]}{x-y} . \tag{24}
\end{equation*}
$$

There have been a lot of literature about the second Gautschi-Kershaw's inequality, please refer to [1, 8, 10, 15, 20, 23, 33, [43]-48].

Our Theorem [2] establishes new bounds, and the series representation for the ratio $\frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x}$.

Theorem 2. Let $x, y>0$ with $x \neq y$, then

$$
\begin{align*}
\frac{1}{3} A(\psi(x), \psi(y)) & +\frac{2}{3} \psi(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(4)}(\max (x, y))<\frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x} \\
& <\frac{1}{3} A(\psi(x), \psi(y))+\frac{2}{3} \psi(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(4)}(\min (x, y)) \tag{25}
\end{align*}
$$

$$
\begin{equation*}
\frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x}=\sum_{k=0}^{\infty} \frac{1}{(2 k+1)!}\left(\frac{y-x}{2}\right)^{2 k} \psi^{(2 k)}\left(\frac{x+y}{2}\right) \tag{26}
\end{equation*}
$$

Proof. It is known that

$$
\begin{equation*}
\psi^{(n)}(x)=(-1)^{n+1} \int_{0}^{\infty} \frac{t^{n}}{1-e^{-t}} e^{-x t} \mathrm{~d} t, \quad n=1,2, \ldots \tag{27}
\end{equation*}
$$

By (12) and (27), we get

$$
\begin{aligned}
& \frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x}=\frac{1}{y-x} \int_{x}^{y} \psi(t) d t \\
= & \frac{1}{6}\left[\psi(x)+4 \psi\left(\frac{x+y}{2}\right)+\psi(y)\right]-\frac{(y-x)^{4}}{2880} \psi^{(4)}(\xi) \\
= & \frac{1}{3} A(\psi(x), \psi(y))+\frac{2}{3} \psi(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(4)}(\xi)
\end{aligned}
$$

for some $\xi$ between $x$ and $y$. This yields (25).
By applying the following result (see [37):

$$
\begin{equation*}
\frac{1}{y-x} \int_{x}^{y} f(t) d t=\sum_{k=0}^{\infty} \frac{1}{(2 k+1)!}\left(\frac{y-x}{2}\right)^{2 k} f^{(2 k)}\left(\frac{x+y}{2}\right) \tag{28}
\end{equation*}
$$

we obtain

$$
\begin{aligned}
& \frac{\ln \Gamma(y)-\ln \Gamma(x)}{y-x}=\frac{1}{y-x} \int_{x}^{y} \psi(t) d t \\
= & \sum_{k=0}^{\infty} \frac{1}{(2 k+1)!}\left(\frac{y-x}{2}\right)^{2 k} \psi^{(2 k)}\left(\frac{x+y}{2}\right) .
\end{aligned}
$$

The proof of Theorem 2 is complete.
Remark 2. Since $\psi^{(2 k)}(x)<0, k=1,2, \ldots$, by (26) we can obtain the second inequality in (21), replacing $y$ and $x$ with $x+t$ and $x+s$, respectively .

The following Theorem 3 presents an extended form of (25).

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Theorem 3. Let $x, y>0$ with $x \neq y$, then for all integers $n \geq 0$,

$$
\begin{align*}
& (-1)^{n}\left[\frac{1}{3} A\left(\psi^{(n+1)}(x), \psi^{(n+1)}(y)\right)+\frac{2}{3} \psi^{(n+1)}(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\min (x, y))\right] \\
< & \frac{(-1)^{n}\left[\psi^{(n)}(y)-\psi^{(n)}(x)\right]}{y-x} \\
< & (-1)^{n}\left[\frac{1}{3} A\left(\psi^{(n+1)}(x), \psi^{(n+1)}(y)\right)+\frac{2}{3} \psi^{(n+1)}(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\max (x, y))\right], \tag{29}
\end{align*}
$$

$$
\begin{equation*}
\frac{\psi^{(n)}(y)-\psi^{(n)}(x)}{y-x}=\sum_{k=0}^{\infty} \frac{1}{(2 k+1)!}\left(\frac{y-x}{2}\right)^{2 k} \psi^{(2 k+n+1)}\left(\frac{x+y}{2}\right) \tag{30}
\end{equation*}
$$

Proof. By (12) and (27), we get

$$
\begin{aligned}
& \frac{\psi^{(n)}(y)-\psi^{(n)}(x)}{y-x}=\frac{1}{y-x} \int_{x}^{y} \psi^{(n+1)}(t) d t \\
= & \frac{1}{6}\left[\psi^{(n+1)}(x)+4 \psi^{(n+1)}\left(\frac{x+y}{2}\right)+\psi^{(n+1)}(y)\right]-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\xi) \\
= & \frac{1}{3} A\left(\psi^{(n+1)}(x), \psi^{(n+1)}(y)\right)+\frac{2}{3} \psi^{(n+1)}(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\xi)
\end{aligned}
$$

for some $\xi$ between $x$ and $y$.
It is easy to see that if $n$ is an even number, then

$$
\begin{align*}
& \frac{1}{3} A\left(\psi^{(n+1)}(x), \psi^{(n+1)}(y)\right)+\frac{2}{3} \psi^{(n+1)}(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\min (x, y)) \\
< & \frac{\psi^{(n)}(y)-\psi^{(n)}(x)}{y-x} \\
< & \frac{1}{3} A\left(\psi^{(n+1)}(x), \psi^{(n+1)}(y)\right)+\frac{2}{3} \psi^{(n+1)}(A(x, y))-\frac{(y-x)^{4}}{2880} \psi^{(n+5)}(\max (x, y)), \tag{31}
\end{align*}
$$

while the reverse inequality holds if $n$ is an odd number. Hence, (29) is valid for all integers $n \geq 0$.

By applying (28), we obtain

$$
\begin{aligned}
& \frac{\psi^{(n)}(y)-\psi^{(n)}(x)}{y-x}=\frac{1}{y-x} \int_{x}^{y} \psi^{(n+1)}(t) d t \\
= & \sum_{k=0}^{\infty} \frac{1}{(2 k+1)!}\left(\frac{y-x}{2}\right)^{2 k} \psi^{(2 k+n+1)}\left(\frac{x+y}{2}\right) .
\end{aligned}
$$

The proof of Theorem 3 is complete.
Remark 3. Since $\psi^{(2 k-1)}(x)>0$ and $\psi^{(2 k)}(x)<0, k=1,2, \ldots$, by (30) we can obtain the first inequality in (23).

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